Quantum machine learning

T. Garipov, P. Izmailov

May 15, 2015

T. Garipov, P. Izmailov

Quantum machine learning

May 15, 2015 1 / 42

1 A brief introduction to quantum mechanics

2 A brief introduction to quantum computations

3 Quantum algorithms for supervised and unsupervised machine learning

1 A brief introduction to quantum mechanics

2 A brief introduction to quantum computations

3 Quantum algorithms for supervised and unsupervised machine learning

D A brief introduction to quantum mechanics

2 A brief introduction to quantum computations

Quantum algorithms for supervised and unsupervised machine learning

D A brief introduction to quantum mechanics

2 A brief introduction to quantum computations

Quantum algorithms for supervised and unsupervised machine learning

$\mathsf{Section}\ 1$

A brief introduction to quantum mechanics

T. Garipov, P. Izmailov

Quantum machine learning

May 15, 2015 3 / 42

э

Motivation

Consider the following experiment

Experiment



- *S* is a source of electrons;
- B is a screen with two holes in it;
- C is another screen, with detectors covering it's surface.

A B A A B A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A

Now consider the plot of the probability of the electron to reach the screen C at the height x as a function of x.



Experiment shows, that

$$P\neq P_1+P_2,$$

where P is the probability distribution with both holes open (subplot a), and *P_i* is the probability distribution with only *i*-th hole open (subplots b and c).

Now consider the plot of the probability of the electron to reach the screen C at the height x as a function of x.



Experiment shows, that

$$P\neq P_1+P_2,$$

where P is the probability distribution with both holes open (subplot a), and P_i is the probability distribution with only *i*-th hole open (subplots b and c).

Postulate 1

Associated to any isolated physical system is a complex Hilbert space known as the *state space* of the system. The system is completely described by its state vector, which is a unit vector in the system's state space.

We will denote the column vector *a* from the state space as $|a\rangle$ and the corresponding row vector as $\langle a|$. We will also denote the inner product of *a* and *b* as $\langle a||b\rangle = \langle a|b\rangle$.

Postulate 1

Associated to any isolated physical system is a complex Hilbert space known as the *state space* of the system. The system is completely described by its state vector, which is a unit vector in the system's state space.

We will denote the column vector a from the state space as $|a\rangle$ and the corresponding row vector as $\langle a|$. We will also denote the inner product of a and b as $\langle a||b\rangle = \langle a|b\rangle$.

Qubit

The simplest quantum system is a *qubit*. It has a 2-dimensional state space \mathcal{H} with orthonormal basis vectors $|0\rangle$ and $|1\rangle$.

The state of a qubit can be written in the form

 $|\psi
angle = a|0
angle + b|1
angle, \qquad$ where $a,b\in\mathbb{C}$

As a state vector is a unit vector, we have

$$\langle \psi | \psi \rangle = |a|^2 + |b|^2 = 1.$$

Qubit

The simplest quantum system is a *qubit*. It has a 2-dimensional state space \mathcal{H} with orthonormal basis vectors $|0\rangle$ and $|1\rangle$.

The state of a qubit can be written in the form

 $|\psi
angle = a|0
angle + b|1
angle, \qquad$ where $a,b\in\mathbb{C}$

As a state vector is a unit vector, we have

$$\langle \psi | \psi \rangle = |\mathbf{a}|^2 + |\mathbf{b}|^2 = 1.$$

Evolution

Postulate 2

The evolution of a closed quantum system is described by a unitary transformation. That is, the state $|\psi\rangle$ of the system at time t_1 is related to the state $|\psi'\rangle$ of the system at time t_2 by a unitary operator U which depends only on the times t_1 and t_2 .

Postulate 2'

The evolution of the state of a closed quantum system is described by the Schrödinger equation,

$$i\hbar \frac{d}{dt}|\psi(t)
angle = H(t)|\psi(t)
angle,$$

where \hbar is the Planks constant and H is a hermitian operator (that might depend on time) known as the *Hamiltonian* of a closed system.

Evolution

Postulate 2

The evolution of a closed quantum system is described by a unitary transformation. That is, the state $|\psi\rangle$ of the system at time t_1 is related to the state $|\psi'\rangle$ of the system at time t_2 by a unitary operator U which depends only on the times t_1 and t_2 .

Postulate 2'

The evolution of the state of a closed quantum system is described by the Schrödinger equation,

$$i\hbar rac{d}{dt} |\psi(t)
angle = H(t)|\psi(t)
angle,$$

where \hbar is the Planks constant and H is a hermitian operator (that might depend on time) known as the *Hamiltonian* of a closed system.

Consider the following example of the quantum system evolution.

Hadamard gate

The matrix representation of the Hadamard gate is the following

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1\\ 1 & -1 \end{bmatrix}$$

It's easy to see, that

$$\begin{split} H|0\rangle &= (|0\rangle + |1\rangle)/\sqrt{2}, \\ H|1\rangle &= (|0\rangle - |1\rangle)/\sqrt{2}. \end{split}$$

Postulate 3

Quantum measurements are described by a collection $\{M_m\}$ of *measurment operators*, acting on the state space of the system being measured. The index *m* referes to the measurement outcomes. If the system is in the state $|\psi\rangle$ immediately before the measurement then the probability that *m* occurs is

$$p(m) = \langle \psi | M_m^* M_m | \psi \rangle,$$

and the state after the measurement is

$$rac{M_m|\psi
angle}{\sqrt{\langle\psi|M_m^*M_m|\psi
angle}}.$$

As an example, consider *the measurement of a qubit in the computational basis*.

The measurement of a qubit in the computational basis

This is a measurement on a single qubit with two outcomes, defined by the operators

$$M_0 = |0
angle\langle 0|,$$

 $M_1 = |1
angle\langle 1|.$

If the state of the system before the measurement is $|\psi\rangle = a|0\rangle + b|1\rangle$, then the probability to obtain a 0 outcome, for example, is

$$\langle \psi | \mathcal{M}_0^* \mathcal{M}_0 | \psi \rangle = \langle \psi | \mathcal{M}_0 | \psi \rangle = |\mathbf{a}|^2.$$

Postulate 4

The state space of a composite system is a tensor product of the state spaces of the component physical systems. Moreover, if we have systems numbered 1 to n, and *i*-th system is in the state $|\psi_i\rangle$, then the joint state of the total system is $|\psi_1\rangle \otimes \ldots \otimes |\psi_n\rangle$.

State space of a 2 qubit system

For example, for two qubits the state vector is of the form

$$|\psi
angle = a|00
angle + b|01
angle + c|10
angle + d|11
angle,$$

where $|a|^2 + |b|^2 + |c|^2 + |d|^2 = 1$.

Postulate 4

The state space of a composite system is a tensor product of the state spaces of the component physical systems. Moreover, if we have systems numbered 1 to n, and *i*-th system is in the state $|\psi_i\rangle$, then the joint state of the total system is $|\psi_1\rangle \otimes \ldots \otimes |\psi_n\rangle$.

State space of a 2 qubit system

For example, for two qubits the state vector is of the form

$$|\psi\rangle = a|00\rangle + b|01\rangle + c|10\rangle + d|11\rangle,$$

where $|a|^2 + |b|^2 + |c|^2 + |d|^2 = 1$.

Entagled state is a quantum state $|\Psi\rangle$ of S_1 , S_2 system that can't be represented as a tensor product of any states $|\Psi_{S_1}\rangle$ and $|\Psi_{S_2}\rangle$ from S_1 and S_2 state spaces respectively.

 $|\Psi_{\mathcal{S}_1}\rangle\otimes|\Psi_{\mathcal{S}_2}\rangle\neq|\Psi\rangle$

Entaglement example

Consider a two-qubit state $|\phi\rangle=\frac{|00
angle+|11
angle}{\sqrt{2}}$. It is easy to show that this is an entagled state.

Section 2

A brief introduction to quantum computations

T. Garipov, P. Izmailov

Quantum machine learning

May 15, 2015 14 / 42

э

Main steps of qunatum computation:

- Initialization quantum states.
- 2 Performance of unitary operations on some subsets of qubits.
- State measurement of the interesting set of qubits.

Almost always correct outcome of computations is obtained with some probability or algorithm aims to estimate some probability.

Quantum computer includes two parts.

- Quantum part:
 - Initialize quantum memory.
 - Perform an unitary operation on some subset of qubits.
 - Make a measurement of a quantum state.
- Classical part:
 - Choose the current unitary operation and a subset of qubits to perform it on.
 - Controll the computation process.

- State of a set of N qubits is a superposition of 2^N possible states.
- The result of a measurement is one of 2^N states, which is obtained with some probability.
- So the quantum memory is principially more powerfull than classical memory.
- One quantum operation makes much more work than one classiscal computer operation.
- Main source of speed-up is the existence of entagled states.

• Add an ancillary qubit in the state $|0\rangle$.

Apply the Hadamard transform to the ancillary qubit.

$$H = \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix}$$

- In the controlled swap procedure (CSWAP).
- O Apply another Hadamard transform to the ancillary qubit.
- Measure state of the ancillary qubit. If $|\phi\rangle = |\psi\rangle$ then $|0\rangle$ is obtained with probability 1.

- Add an ancillary qubit in the state $|0\rangle$.
- Output the Hadamard transform to the ancillary qubit.

$$H = \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix}$$

- In the controlled swap procedure (CSWAP).
- Apply another Hadamard transform to the ancillary qubit.
- Measure state of the ancillary qubit. If $|\phi\rangle = |\psi\rangle$ then $|0\rangle$ is obtained with probability 1.

- Add an ancillary qubit in the state $|0\rangle$.
- Output the Hadamard transform to the ancillary qubit.

$$H = \left[\begin{array}{cc} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{array} \right]$$

- Series Perform the controlled swap procedure (CSWAP).
- Apply another Hadamard transform to the ancillary qubit.
- Measure state of the ancillary qubit. If $|\phi\rangle = |\psi\rangle$ then $|0\rangle$ is obtained with probability 1.

- Add an ancillary qubit in the state $|0\rangle$.
- Output the Hadamard transform to the ancillary qubit.

$$H = \left[\begin{array}{cc} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{array} \right]$$

- Seriorm the controlled swap procedure (CSWAP).
- Opply another Hadamard transform to the ancillary qubit.
- Measure state of the ancillary qubit. If $|\phi\rangle = |\psi\rangle$ then $|0\rangle$ is obtained with probability 1.

- Add an ancillary qubit in the state $|0\rangle$.
- Output the Hadamard transform to the ancillary qubit.

$$H = \left[\begin{array}{cc} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{array} \right]$$

- Series of the controlled swap procedure (CSWAP).
- O Apply another Hadamard transform to the ancillary qubit.
- Measure state of the ancillary qubit. If $|\phi\rangle = |\psi\rangle$ then $|0\rangle$ is obtained with probability 1.

The evolution through algorithm steps:

$$\begin{split} |0\rangle|\phi\rangle|\psi\rangle &\xrightarrow{H} \frac{|0\rangle+|1\rangle}{\sqrt{2}} |\phi\rangle|\psi\rangle \xrightarrow{CSWAP} \frac{|0\rangle|\phi\rangle|\psi\rangle+|1\rangle|\psi\rangle|\phi\rangle}{\sqrt{2}} \xrightarrow{H} \\ &\frac{|0\rangle\left[|\phi\rangle|\psi\rangle+|\psi\rangle|\phi\rangle\right]+|1\rangle\left[|\phi\rangle|\psi\rangle-|\psi\rangle|\phi\rangle\right]}{2} \end{split}$$

The probability of passing the test:

$${{P}=rac{1}{4}(\langle \phi | \langle \psi |+ \langle \psi | \langle \phi |) (| \phi
angle | \psi
angle + | \psi
angle | \phi
angle)=rac{1+|\langle \phi | \psi
angle |^2}{2}}{2}}$$

T. Garipov, P. Izmailov

Grover's search algorithm

Task: Find a solution of a euqation f(x) = 1, where $x \in \{0, 1\}^n$, and f(x) is a boolean function. Time: $O(\sqrt{2^N})$ Memory: O(N)

Shor's factorization algorithm

Task: Factorization of number N.Time: $O([\log N]^3)$ Goodbye RSA!Memory: $O(\log N)$

Quantum algorithm for linear systems of equations

Task: Solving a linear system of equations with N variables. **Assumption:** The system is sparse and has a low conditional number κ . **Time:** $O(\kappa \log N)$ (versus classical $O(\kappa N)$) Theory of quantum computations is quite well developed and it have shown possible huge potential of quantum computers compared to classical computers.

But modern implementations of quantum computer are strongly restricted and are used only for specialized tasks. Building a full quantum computer as a real physical device is a fundamental problem of 21th century physics.

Quantum computations and machine learning

Motivation of quantum machine learning:

- There is a challenge that machine learning with rapidly growing "big data" could become intractable for classical computers.
- Manipulation with high-dimensional vector is a core routine for machine learnig algorithms.
- Quantum computers appears to be good at such manipulations.

Vector storing

Stroring a representation of a 2^N -dimensional unit vector require only N qubits. And such representations could be constructed in O(N) time.

• Quantum machine learning may provide an exponential speed-up over known algorithms for problems involving evaluating distances and inner products between large vectors.

Section 3

Quantum algorithms for supervised and unsupervised machine learning

T. Garipov, P. Izmailov

Quantum machine learning

May 15, 2015 23 / 42

We assume, that the data sets that consist of vectors and collections of vectors are originally stored in qRAM.

Then constructing the $\log_2 N$ qubit quantum state $|v\rangle = |v|^{-1/2}v$ takes $O(\log_2 N)$ steps.
Consider the task of assigning a post-processed vector $u \in \mathbb{R}^n$ to one of two sets V, W given M representatives of each classes.

A common method for such an assignment is evaluating the distances

$$u - \frac{1}{M} \sum_{j=1}^{M} v_j$$

and

$$u-\frac{1}{M}\sum_{j=1}^{M}w_{j}\bigg|,$$

and assign the vector to the class, for which this distance is smaller.

distance determining algorithm

• Create an ancillary variable with M + 1 states.

Construct the state

$$|\psi
angle = rac{1}{\sqrt{2}}(|0
angle|u
angle + rac{1}{\sqrt{M}}\sum_{j=1}^{M}|j
angle|v_{j}
angle).$$

Use the swap-test to determine, weather the ancillary variable is in the state

$$ert arphi
angle = rac{1}{\sqrt{Z}} (ert u ert ert 0
angle - rac{1}{\sqrt{M}} \sum_{j=1}^M ert v_j ert ert j
angle),$$

where
$$Z = |u|^2 + (1/M) \sum_{j=1}^M |v_j|^2.$$

The probability of success in this measurement is equal to the distance $|u - (1/M) \sum_{i=1}^{M} v_i|^2$ divided by Z. Thus, performing this measurement

several times we can estimate the distance

T. Garipov, P. Izmailov

Quantum machine learning

distance determining algorithm

• Create an ancillary variable with M + 1 states.

Onstruct the state

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle|u\rangle + \frac{1}{\sqrt{M}}\sum_{j=1}^{M}|j\rangle|v_{j}\rangle).$$

Use the swap-test to determine, weather the ancillary variable is in the state

$$ert arphi
angle = rac{1}{\sqrt{Z}} (ert u ert ert 0
angle - rac{1}{\sqrt{M}} \sum_{j=1}^M ert v_j ert ert j
angle),$$

where
$$Z=|u|^2+(1/M)\sum\limits_{j=1}^M|v_j|^2.$$

The probability of success in this measurement is equal to the distance $|u - (1/M) \sum_{i=1}^{M} v_i|^2$ divided by Z. Thus, performing this measurement

several times we can estimate the distance

T. Garipov, P. Izmailov

Quantum machine learning

distance determining algorithm

- Create an ancillary variable with M + 1 states.
- Onstruct the state

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle|u\rangle + \frac{1}{\sqrt{M}}\sum_{j=1}^{M}|j\rangle|v_{j}\rangle).$$

Use the swap-test to determine, weather the ancillary variable is in the state

$$|arphi
angle = rac{1}{\sqrt{Z}}(|u||0
angle - rac{1}{\sqrt{M}}\sum_{j=1}^{M}|v_j||j
angle),$$

where
$$Z = |u|^2 + (1/M) \sum_{j=1}^M |v_j|^2.$$

The probability of success in this measurement is equal to the distance $|u - (1/M) \sum_{i=1}^{M} v_i|^2$ divided by Z. Thus, performing this measurement

several times we can estimate the distance

T. Garipov, P. Izmailov

Quantum machine learning

distance determining algorithm

- Create an ancillary variable with M + 1 states.
- Onstruct the state

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle|u\rangle + \frac{1}{\sqrt{M}}\sum_{j=1}^{M}|j\rangle|v_{j}\rangle).$$

Use the swap-test to determine, weather the ancillary variable is in the state

$$|arphi
angle = rac{1}{\sqrt{Z}}(|u||0
angle - rac{1}{\sqrt{M}}\sum_{j=1}^{M}|v_j||j
angle),$$

where
$$Z = |u|^2 + (1/M) \sum_{j=1}^M |v_j|^2.$$

The probability of success in this measurement is equal to the distance $|u - (1/M) \sum_{j=1}^{M} v_j|^2$ divided by Z. Thus, performing this measurement several times we can estimate the distance

T. Garipov, P. Izmailov

Quantum machine learning

As already been mentioned, the time evolution of a quantum system is described by the Schrödinger equation:

$$i\hbarrac{d}{dt}|\psi(t)
angle=H(t)|\psi(t)
angle,$$

where H(t) is the Hamiltonian of the system. A Hamiltonian is described by a hermitian matrix, whose eigenvectors represent the *eigenstates* of the system. The corresponding eigenvalues refer to the different energies of the eigenstates. The state with the lowest energy is called the *ground state* of the system.

Quantum adiabatic theorem

A physical system that is initially in its ground state, tends to stay in this lowest energy state, provided that the Hamiltonian of the system is changed 'slowly enough'. As already been mentioned, the time evolution of a quantum system is described by the Schrödinger equation:

$$i\hbarrac{d}{dt}|\psi(t)
angle=H(t)|\psi(t)
angle,$$

where H(t) is the Hamiltonian of the system. A Hamiltonian is described by a hermitian matrix, whose eigenvectors represent the *eigenstates* of the system. The corresponding eigenvalues refer to the different energies of the eigenstates. The state with the lowest energy is called the *ground state* of the system.

Quantum adiabatic theorem

A physical system that is initially in its ground state, tends to stay in this lowest energy state, provided that the Hamiltonian of the system is changed 'slowly enough'.

A B A B A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A

Adiabatic quantum computations framework

- Define the Hamiltonian H_0 so that it's eigenstates are easy to compute. Set the system to the ground state of H_0 .
- 2 Define the final Hamiltonian as

$$H_f = \sum_{z \in \{0,1\}^n} f(z) |z\rangle \langle z|.$$

Gradually transform the Hamiltonian of the system from H₀ to H_f.
 The final state of the system is the minimizer of the objective function.

Adiabatic quantum computations framework

- Define the Hamiltonian H_0 so that it's eigenstates are easy to compute. Set the system to the ground state of H_0 .
- 2 Define the final Hamiltonian as

$$H_f = \sum_{z \in \{0,1\}^n} f(z) |z\rangle \langle z|.$$

Gradually transform the Hamiltonian of the system from H₀ to H_f.
 The final state of the system is the minimizer of the objective function.

Adiabatic quantum computations framework

- Define the Hamiltonian H_0 so that it's eigenstates are easy to compute. Set the system to the ground state of H_0 .
- 2 Define the final Hamiltonian as

$$H_f = \sum_{z \in \{0,1\}^n} f(z) |z\rangle \langle z|.$$

Gradually transform the Hamiltonian of the system from H₀ to H_f.
 The final state of the system is the minimizer of the objective function

Adiabatic quantum computations framework

- Define the Hamiltonian H_0 so that it's eigenstates are easy to compute. Set the system to the ground state of H_0 .
- 2 Define the final Hamiltonian as

$$H_f = \sum_{z \in \{0,1\}^n} f(z) |z\rangle \langle z|.$$

Gradually transform the Hamiltonian of the system from H₀ to H_f.
The final state of the system is the minimizer of the objective function.

The standard classical algorithm for doing this is the Lloyd's algorithm:

Lloyd algorithm

- Choose k initial centroids.
- 2 Assign each vector to the cluster with the closest mean.
- ③ Recalculate the centroids of the clusters.
- Repeat steps (2–3) until a stationary assignment is attained.

The standard classical algorithm for doing this is the Lloyd's algorithm:

Lloyd algorithm

- Choose k initial centroids.
- 2 Assign each vector to the cluster with the closest mean.
- ③ Recalculate the centroids of the clusters.
- Repeat steps (2–3) until a stationary assignment is attained.

The standard classical algorithm for doing this is the Lloyd's algorithm:

Lloyd algorithm

- Choose k initial centroids.
- 2 Assign each vector to the cluster with the closest mean.
- **③** Recalculate the centroids of the clusters.

Repeat steps (2–3) until a stationary assignment is attained.

The standard classical algorithm for doing this is the Lloyd's algorithm:

Lloyd algorithm

- Choose k initial centroids.
- 2 Assign each vector to the cluster with the closest mean.
- **③** Recalculate the centroids of the clusters.
- Repeat steps (2-3) until a stationary assignment is attained.

- Choose k initial centroids i_1, \ldots, i_k .
- **②** Using the adiabatic quantum computations find the re-clustering that minimizes the distances from each point to the corresponding centroid. The result is a state $|\psi_1\rangle = \frac{1}{\sqrt{M}} \sum |c\rangle |j\rangle$
- Construct multiple copies of this state and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_1^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Use the adiabatic quantum computations to find the re-clustering that minimizes the distances from each point to the corresponding centroid of the clasters, obtained on the previous iteration. The result is a state |\u03c6_i\u03c6_i.
- Construct multiple copies of the state $|\psi_i\rangle$ and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_i^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Repeat the steps (4 5) until two successive states |\u03c6_i\u03c6 and |\u03c6_{i+1}\u03c6_i coincide, which can be verified using the swap test.

- Choose k initial centroids i_1, \ldots, i_k .
- **2** Using the adiabatic quantum computations find the re-clustering that minimizes the distances from each point to the corresponding centroid. The result is a state $|\psi_1\rangle = \frac{1}{\sqrt{M}} \sum_{c \ i \in c} |c\rangle |j\rangle$
- (a) Construct multiple copies of this state and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_1^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Use the adiabatic quantum computations to find the re-clustering that minimizes the distances from each point to the corresponding centroid of the clasters, obtained on the previous iteration. The result is a state |\u03c6_i\u03c6_i.
- Onstruct multiple copies of the state |ψ_i⟩ and perform projective measurements of the |c⟩. Thus we obtain individual cluster states |φ^c_i⟩ = 1/√M ∑_{i∈c} |j⟩.
- Repeat the steps (4 5) until two successive states $|\psi_i\rangle$ and $|\psi_{i+1}\rangle$ coincide, which can be verified using the swap test.

- Choose k initial centroids i_1, \ldots, i_k .
- **2** Using the adiabatic quantum computations find the re-clustering that minimizes the distances from each point to the corresponding centroid. The result is a state $|\psi_1\rangle = \frac{1}{\sqrt{M}} \sum_{c \ i \in c} |c\rangle |j\rangle$
- Construct multiple copies of this state and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_1^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Use the adiabatic quantum computations to find the re-clustering that minimizes the distances from each point to the corresponding centroid of the clasters, obtained on the previous iteration. The result is a state |\u03c6_i\u03c6.
- Construct multiple copies of the state $|\psi_i\rangle$ and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_i^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Repeat the steps (4 5) until two successive states |\u03c6_i\u03c6 and |\u03c6_{i+1}\u03c6_i coincide, which can be verified using the swap test.

- Choose k initial centroids i_1, \ldots, i_k .
- **2** Using the adiabatic quantum computations find the re-clustering that minimizes the distances from each point to the corresponding centroid. The result is a state $|\psi_1\rangle = \frac{1}{\sqrt{M}} \sum_{c \ i \in c} |c\rangle |j\rangle$
- Construct multiple copies of this state and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_1^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Use the adiabatic quantum computations to find the re-clustering that minimizes the distances from each point to the corresponding centroid of the clasters, obtained on the previous iteration. The result is a state |\u03c6_i\u03c6.
- Construct multiple copies of the state $|\psi_i\rangle$ and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_i^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Seperat the steps (4 5) until two successive states |\u03c6_i\u03c6 and |\u03c6_{i+1}\u03c6_i coincide, which can be verified using the swap test.

- Choose k initial centroids i_1, \ldots, i_k .
- **2** Using the adiabatic quantum computations find the re-clustering that minimizes the distances from each point to the corresponding centroid. The result is a state $|\psi_1\rangle = \frac{1}{\sqrt{M}} \sum_{c \ i \in c} |c\rangle |j\rangle$
- Construct multiple copies of this state and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_1^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Use the adiabatic quantum computations to find the re-clustering that minimizes the distances from each point to the corresponding centroid of the clasters, obtained on the previous iteration. The result is a state |\u03c6_i\u03c6.
- Construct multiple copies of the state $|\psi_i\rangle$ and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_i^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- ⁽³⁾ Repeat the steps (4-5) until two successive states $|\psi_i\rangle$ and $|\psi_{i+1}\rangle$ coincide, which can be verified using the swap test.

- Choose k initial centroids i_1, \ldots, i_k .
- **2** Using the adiabatic quantum computations find the re-clustering that minimizes the distances from each point to the corresponding centroid. The result is a state $|\psi_1\rangle = \frac{1}{\sqrt{M}} \sum_{c \ i \in c} |c\rangle |j\rangle$
- Construct multiple copies of this state and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_1^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Use the adiabatic quantum computations to find the re-clustering that minimizes the distances from each point to the corresponding centroid of the clasters, obtained on the previous iteration. The result is a state |\u03c6_i\u03c6.
- Construct multiple copies of the state $|\psi_i\rangle$ and perform projective measurements of the $|c\rangle$. Thus we obtain individual cluster states $|\varphi_i^c\rangle = \frac{1}{\sqrt{M}} \sum_{i \in c} |j\rangle$.
- Repeat the steps (4 5) until two successive states $|\psi_i\rangle$ and $|\psi_{i+1}\rangle$ coincide, which can be verified using the swap test.

The output of the algorithm is a state

$$|\chi\rangle = rac{1}{\sqrt{M}}\sum_{j}|c_{j}
angle |j
angle$$

that contains the labels j of vectors, correlated with their cluster assignments c_j . Sampling from this state we obtain a statistical picture of the clustering.

For the vector assignment algorithm the computational complexity estimate is $O(\log MN)$.

For the quantum Lloyd algorithm the computational complexity of constructing the state $|\chi\rangle$ estimation is $O(k \log kMN)$.

For the vector assignment algorithm the computational complexity estimate is $O(\log MN)$.

For the quantum Lloyd algorithm the computational complexity of constructing the state $|\chi\rangle$ estimation is $O(k \log kMN)$.

Section 4

Algorithm implementation and experiments

T. Garipov, P. Izmailov

Quantum machine learning

May 15, 2015 33 / 42

æ

Task: Classification of *N*-dimensional vectors (where N = 2, 4, 8) to two classes with one reference vector in each class.

Notation:

- Classes are labeled as A and B.
- \vec{v}_A and \vec{v}_B are reference vectors.
- \vec{u} is the new sample vector.

Classification of the new sample is done by comparing distances:

$$D_A = |\vec{u} - \vec{v}_A|, \qquad D_B = |\vec{u} - \vec{v}_B|$$

Compting distance $|\vec{u} - \vec{v}|$:

Represent vectors with quantum states:

$$\vec{u} = |u||u\rangle$$
, $\vec{v} = |v||v\rangle$

2 Add an ancillary qubit and create an entagled state:

$$\left|\phi
ight
angle=(\left|0
ight
angle_{\mathit{anc}}\left|u
ight
angle_{\mathit{new}}+\left|1
ight
angle_{\mathit{anc}}\left|v
ight
angle_{\mathit{ref}})/\sqrt{2}$$

$$|\psi
angle = (|u||0
angle - |v||1
angle)/\sqrt{|u|^2 + |v|^2}$$

- Estimate the succes probability of measurement p by repeated measurements.
- Solution $|\vec{u} \vec{v}|$ using p.

Compting distance $|\vec{u} - \vec{v}|$:

Represent vectors with quantum states:

$$\vec{u} = |u||u\rangle, \qquad \vec{v} = |v||v\rangle$$

Add an ancillary qubit and create an entagled state:

$$\left|\phi
ight
angle=(\left|0
ight
angle_{\mathit{anc}}\left|u
ight
angle_{\mathit{new}}+\left|1
ight
angle_{\mathit{anc}}\left|v
ight
angle_{\mathit{ref}})/\sqrt{2}$$

$$|\psi\rangle = (|u||0\rangle - |v||1\rangle)/\sqrt{|u|^2 + |v|^2}$$

- Estimate the succes probability of measurement p by repeated measurements.
- Solution $|\vec{u} \vec{v}|$ using p.

Compting distance $|\vec{u} - \vec{v}|$:

Represent vectors with quantum states:

$$\vec{u} = |u||u\rangle, \qquad \vec{v} = |v||v\rangle$$

Add an ancillary qubit and create an entagled state:

$$\ket{\phi} = (\ket{0}_{\mathit{anc}} \ket{u}_{\mathit{new}} + \ket{1}_{\mathit{anc}} \ket{v}_{\mathit{ref}})/\sqrt{2}$$

$$|\psi\rangle = (|u||0\rangle - |v||1\rangle)/\sqrt{|u|^2 + |v|^2}$$

- Estimate the succes probability of measurement p by repeated measurements.
- Solution $|\vec{u} \vec{v}|$ using p.

Compting distance $|\vec{u} - \vec{v}|$:

Represent vectors with quantum states:

$$\vec{u} = |u||u\rangle, \qquad \vec{v} = |v||v\rangle$$

Add an ancillary qubit and create an entagled state:

$$\ket{\phi} = (\ket{0}_{\mathit{anc}} \ket{u}_{\mathit{new}} + \ket{1}_{\mathit{anc}} \ket{v}_{\mathit{ref}})/\sqrt{2}$$

Make a measurement on the ancillary qubit, projecting it onto the state:

$$|\psi
angle = (|u||0
angle - |v||1
angle)/\sqrt{|u|^2 + |v|^2}$$

Stimate the succes probability of measurement p by repeated measurements.

Compting distance $|\vec{u} - \vec{v}|$:

Represent vectors with quantum states:

$$\vec{u} = |u||u\rangle, \qquad \vec{v} = |v||v\rangle$$

Add an ancillary qubit and create an entagled state:

$$\ket{\phi} = (\ket{0}_{\mathit{anc}} \ket{u}_{\mathit{new}} + \ket{1}_{\mathit{anc}} \ket{v}_{\mathit{ref}})/\sqrt{2}$$

$$|\psi
angle = (|u||0
angle - |v||1
angle)/\sqrt{|u|^2 + |v|^2}$$

- Stimate the succes probability of measurement p by repeated measurements.
- $o Calculate |\vec{u} \vec{v}| using p.$

The distance between \vec{u} and \vec{v} can be directly calculated from p:

$$|\vec{u} - \vec{v}| = \sqrt{2p(|u|^2 + |v|^2)}$$

Also the inner product between $|u\rangle$ and $|v\rangle$ can be obtained:

$$\langle u|v\rangle = (0.5 - p)(|u|^2 + |v|^2)/(|u||v|)$$

Experiments: setup



Figure: Experimental setup with 4 photonic qubits

T. Garipov, P. Izmailov

Quantum machine learning

 $\exists \rightarrow b$ May 15, 2015 37 / 42

э

Experiments: results 2D classification



Experiment for 2D vectors:

- Only 2 of 100 samples are misclassified.
- Errors occured at vectors that are close to boundary where absolute error of probability estimation is close to $|D_A - D_B|$.
- Estimation of probability done by 10 000 repeated measurements for each vector.
- Time needed for processing one vector is 1 sec.

Experiments: results 4D classification

		D _A -D _B Theory Exp:		Group	Correct?
1	(2.00, 0.00, 0.00, 0.00)	-1.45	-0.93	Α	/
2	(0.00, 0.00, 0.00, 2.00)	0.82	0.50	В	1
3	(0.35, 0.20, 0.00, 0.00)	-0.79	-0.71	А	1
4	(0.23, 0.19, 0.08, 0.07)	-0.54	-0.51	А	1
5	(1.32, 3.62, 1.57, 4.32)	0.74	0.48	В	1
6	(0.15, 0.17, 0.82, 0.98)	1.26	0.72	В	1
7	(0.18, 0.10, 1.02, 0.59)	0.98	0.76	В	1
8	(0.97, 0.17, 0.17, 0.03)	-1.37	-0.93	Α	1
9	(0.68, 0.25, 0.00, 0.00)	-1.18	-0.79	Α	1
10	(0.83, 0.48, 1.44, 0.83)	0.67	0.17	В	1
11	(1.27, 1.06, 3.48, 2.92)	1.13	0.76	В	1
12	(0.40, 0.40, 0.40, 0.40)	-0.10	-0.26	А	1
13	(0.09, 0.15, 0.49, 0.85)	0.80	0.55	В	1
14	(0.10, 0.55, 0.06, 0.32)	-0.19	-0.28	А	1
15	(1.94, 0.34, 0.34, 0.06)	-1.22	-1.10	А	1
16	(3.42, 1.24, 1.97, 0.72)	-0.34	-0.39	А	1
17	(0.66, 0.00, 1.80, 0.00)	0.40	-0.02	А	х

Experiment for 4D vectors:

• $\vec{v}_A = (1, 0, 0, 0).$

•
$$\vec{v}_B = (0, 0, 1, 1).$$

- 500 measurements per estimation.
- Data acquisition time for one vector is 2 min.

Experiments: results

8D classification

		D _A -D _B Theory Exp:		Group	Correct?
1	(2.00, 0.00, 0.00, 0.00, 0.00, 0.00, 0.00, 0.00)	-1.24	-0.84	A	/
2	(0.00, 0.00, 0.00, 0.00, 0.00, 0.00, 0.00, 0.60)	0.77	0.55	В	1
3	(1.77, 0.00, 0.00, 0.00, 1.24, 0.00, 0.00, 0.00)	-0.92	-0.52	А	1
4	(0.40, 0.23, 0.11, 0.06, 0.03, 0.02, 0.01, 0.01)	-0.45	-0.14	А	1
5	(0.00, 0.00, 1.23, 1.23, 0.00, 0.00, 0.33, 0.33)	0.17	0.10	В	1
6	(0.30, 0.03, 0.30, 0.03, 1.12, 0.10, 1.12, 0.10)	-0.11	-0.24	A	1
7	(0.42, 0.90, 0.35, 0.76, 0.00, 0.00, 0.00, 0.00)	-0.28	-0.21	A	1
8	(0.54, 0.54, 0.00, 0.00, 0.54, 0.54, 0.00, 0.00)	-0.43	-0.50	A	1
9	(0.11, 1.24, 0.19, 2.15, 0.06, 0.72, 0.11, 1.24)	0.40	-0.17	А	×

Experiment for 8D vectors:

- $\vec{v}_A = (1, 0, 0, 0, 0, 0, 0, 0).$
- $\vec{v}_B = (0, 0, 0, 0, 0, 0, 0, 1).$
- 500 measurements per estimation.
- Data acquisition time for one vector is 4 min.
Experiments summary:

- First experimental demonstration of machine learning on a photonic quantum computer.
- Experimental prove of suitability and potential power of quantum machine learning.
- Future studies are planned to design circuits that will be able to achieve speed-up not only with respect to N but also with respect to number of training samples M.

- 🔋 S. Lloyd, M. Mohseni, and P. Rebentrost, arXiv:1307.0411.
- X.-D. Cai, D. Wu, Z.-E. Su, M.-C. Chen, X.-L. Wang, L. Li, N.-L. Liu, C.-Y. Lu, J.-W. Pan, arXiv:1409.7770v3.